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Solar-driven photothermal nanostructured materials designs and prerequisites for evaporation and catalysis applications

Liangliang Zhu,^a Minmin Gao,^a Connor Kang Nuo Peh^a and Ghim Wei Ho^{id} *^{abc}

Solar energy is a major source of renewable energy with the potential to meet the energy demand and to support the sustainable development of the world. The efficient harvesting and conversion of solar energy is one of the key factors to maximize the utilization of solar energy. In general, solar energy can be harnessed and converted into various kinds of energy, including electricity, fuels and thermal energy, through photovoltaic, photochemical and photothermal processes, respectively. Among these technologies, photothermal conversion is a direct conversion process that has attained the highest achievable conversion efficiency. The photothermal effect has been used as a novel strategy to augment vaporization and catalysis performance. In this review, we look into the basis of the photothermal conversion process, the design of efficient photothermal conversion materials in terms of both light harvesting and thermal management, a fundamental understanding of various system schemes, and the recent progress in photothermal evaporation and catalysis applications. This review aims to afford researchers with a better understanding of the photothermal effect and provide a guide for the rational design and development of highly efficient photothermal materials in energy and environmental fields.

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1. Introduction

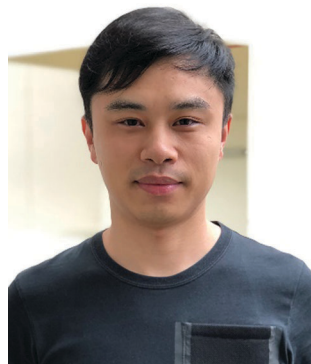
The global energy demand is ever-growing, driven by strong economic growth and an increasing pursuit of a high-quality

of life.^{1,2} With rising concerns about the depletion of non-renewable resources alongside environmental issues, satisfying future energy needs in a sustainable way is a compelling need. Besides, the efficient utilization of renewable solar energy has attracted a lot of attention for decades as one of the most important strategies in green energy innovation. The critical challenge with this energy source is to enhance the utilization efficiency by extending the solar spectrum response from the UV to the visible, and even to the infrared (IR) region.^{3,4} Photothermal conversion is an important pathway to harvest

^a Department of Electrical and Computer Engineering, National University of Singapore, 4 Engineering Drive 3, 117583, Singapore. E-mail: elehgw@nus.edu.sg

^b Engineering Science Programme, National University of Singapore, 9 Engineering Drive 1, 117575, Singapore

^c Institute of Materials Research and Engineering, A*STAR (Agency for Science, Technology and Research), 3 Research Link, 117602, Singapore



Liangliang Zhu

Liangliang Zhu received his PhD degree in Material Chemical Engineering at Nanjing University of Technology, China, in 2013. He is now a research fellow at the Department of Electrical & Computer Engineering at National University of Singapore (NUS). His current research interest is the preparation of photothermal materials for water evaporation and energy generation.



Minmin Gao

Minmin Gao received her BEng (Hons) degree in Electrical Engineering from National University of Singapore (NUS) in 2012, as well as her PhD in 2016 under the supervision of Associate Professor Dr Ghim Wei Ho. Her research interest currently focuses on TiO₂-based core-shell nanostructures for catalytic applications.

solar energy and to enable its use in a broad range of applications, including environmental purification, water distillation/desalination and the catalytic production of fuels and chemicals.^{5–7} Solar water evaporation is highly promising for fresh water production for meeting remote areas/emergency needs, while photothermal catalysis provides an effective method to energetically drive solar fuels catalytic processes. However, the photothermal conversion efficiency is mainly limited by sunlight absorption and thermal management. Accordingly, rational material and system designs that favour efficient photon absorption and a reduction of superfluous heat losses are highly sought after.

Recent efforts in the development of solar absorber nano-materials capable of broadband absorption, *i.e.* the visible to IR radiation, have become increasingly important from the viewpoint of energy conservation and sustainability for photothermal steam generation, water evaporation and catalysis.^{5,8} Herein, we review the recent progress in photothermal evaporation and catalysis applications together with the basis for photothermal conversion and the design of pertinent photothermal converting materials and system configurations for enhanced solar-driven vaporization and catalytic reactivity (Fig. 1). Various solar absorber materials, *e.g.* plasmonic metals,^{9–12} narrow bandgap semiconductors^{13–18} and carbon-based nanostructures,^{19–23} with efficient photothermal conversion capabilities will be discussed. Also, the diverse approaches that have been employed to enhance the light-to-heat conversion efficiency, such as tuning the structural composition, the hybridization of light absorber materials and the optimization of absorber/reactant configurations, will be covered. Finally, conclusions and the future perspectives for photothermal evaporation and catalysis will be presented. The goal of this review is to provide a background into photothermal conversion and to produce a guide that could shed some light on the key points to consider when designing solar absorber nano-materials to aid the better future deployment of solar-driven evaporation and catalysis.

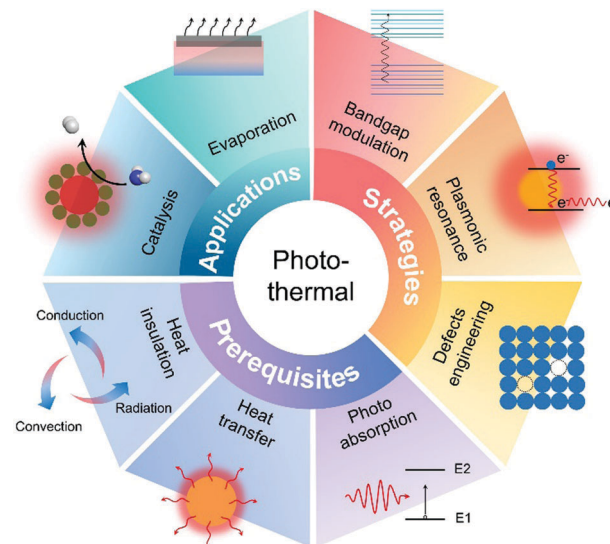


Fig. 1 Schematic diagram of the scope of the present review.

2. Basis of photothermal conversion

It is essential that the optical and thermal properties of a material are considered to achieve an efficient photothermal conversion. The optical properties of materials, such as the absorption coefficient, the matching of a material's absorption spectrum to incident light and heat transfer and loss, play important roles in material selection and in optimization of the light-to-heat conversion (Fig. 2). In addition to the inherent properties of the utilized materials, other factors, such as the thermal stability, mechanical integrity and cost, should also be factored in for practical applications.

2.1. Solar absorptance

Solar absorptance is a measure of the ability of a matter to absorb solar radiation, *i.e.* the ratio of the total absorbed solar radiation to the incident radiation. When a surface of a matter



Connor Kang Nuo Peh

Connor Peh received his BEng (Hons) and PhD degrees from National University of Singapore in 2009 and 2016, respectively. He is currently a research fellow in NUS in A/Prof. Ho Ghim Wei's group. His current research interests primarily focus on clean energy generation, photocatalytic and photoelectrochemical water splitting and the development of materials to harness sunlight in various forms.



Ghim Wei Ho

Ghim Wei Ho received her PhD for work in semiconductor nanostructures at the Nanoscience Centre, University of Cambridge in 2006. She is an associate professor at the Department of Electrical and Computer Engineering at the National University of Singapore (NUS). She is also one of the pioneer associates of the Engineering Science Programme (ESP) at NUS. Her research thrust is focused on the development of functional nanostructured materials for energy, environmental and healthcare applications.

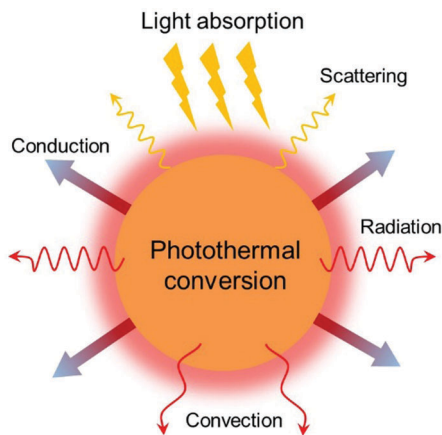


Fig. 2 Schematic diagram of photothermal conversion.

is subjected to electromagnetic radiation, a fraction of the photon energy can be transformed to other forms of energy, such as heat. The total solar absorptance for a given angle of incidence θ is obtained by weighting the spectral absorptance with the solar spectral irradiance distribution of a standard solar spectrum (AM 1.5) and integrating over the wavelengths in which the solar radiation reaches the solar absorber surface. Eqn (2.1) describes the total solar absorptance, $\alpha(\theta)$, of a solar absorber for an incident angle of θ

$$\alpha(\theta) = \frac{\int_{\lambda_{\min}}^{\lambda_{\max}} [1 - R(\theta, \lambda)] A(\theta, \lambda) d\lambda}{\int_{\lambda_{\min}}^{\lambda_{\max}} A(\lambda) d\lambda} \quad (2.1)$$

where λ_{\min} and λ_{\max} are 0.3 μm and 2.5 μm , respectively, and θ is the angle of incidence of light measured from the surface normal of the absorber, $A(\lambda)$ is the wavelength-dependent solar spectral irradiance and $R(\theta, \lambda)$ is the total reflectance at wavelength λ .²⁴

As this review focuses mainly on utilization of the solar spectrum as the source of light (AM 1.5 reference spectrum), the material properties should be directed at achieving the maximum conversion from the sunlight that falls onto the Earth's surface. This means that the material should absorb as much of the electromagnetic radiation as possible from the wavelengths 300 nm to 2500 nm, with the visible range (400–700 nm) comprising $\sim 45\%$ of the power, and the IR region (700–2500 nm) encompassing $\sim 52\%$ of the power. The solar absorber material should be designed to have a high absorption in this region, with negligible reflectance in order to maximize the energy capture from the sun and to convert it to heat energy.

2.2. Heat transfer

In order to improve the efficiencies of solar heat absorbers, attention has to be placed on the heat transfer occurring in the system. Heat flows from one body to another in three main ways: conduction, convection and radiation; driven by thermal gradients between the bodies. Heat conduction occurs when heat flows from a higher temperature source to a lower

temperature body by passing through a medium – passing on the energy by collisions between neighbouring particles.

$$q = kA \frac{(T_1 - T_2)}{L} \quad (2.2)$$

where q is the heat transfer by conduction, k is the thermal conductivity of the material, A is the surface area of heat transfer, T_1 and T_2 are the temperatures of the two different bodies (assuming steady state temperatures) and L is the thickness of the medium.

Heat convection occurs by the motion of a mass of fluid carrying the energy away from the source of heat into a cooler part of the fluid. The convection heat transfer at a constant temperature in a fluid can be expressed as

$$q = \bar{h}(T_{\text{body}} - T_c) \quad (2.3)$$

where q is the heat transfer by convection, \bar{h} is the average heat transfer coefficient, T_{body} is the temperature of the body and T_c is the temperature of the surrounding fluid.

Lastly, heat radiation includes electromagnetic radiation from all objects. No medium of transfer is required and no mass is exchanged in this process. All objects constantly emit radiation based on their temperature and surface properties. Hotter objects usually emit much more radiation than cooler objects, in a relationship proportional to the fourth power of the absolute temperature. All objects therefore emit and absorb radiation at the same time, and the net effect is based on their temperature relative to their surroundings, and their emissivity. The Stefan–Boltzmann law can be re-expressed for a hot object in a cooler environment as

$$P = \varepsilon\sigma A(T_4 - T_c^4) \quad (2.4)$$

where P is the thermal energy emitted per unit area, ε is the emittance of the object, A is the radiating area, T is the temperature of the radiator and T_c is the temperature of the cooler surroundings. Stefan's constant, σ , is taken as $5.6703 \times 10^8 \text{ W m}^{-2} \text{ K}^{-4}$.

As mentioned earlier, heat transfer occurs *via* three ways, and they all depend on a temperature gradient. Conduction is the most likely form of heat transfer for photothermal vaporization and catalysis applications, whereby a short transfer length associated with nanoscale dimensions and a high surface area coupled with an appropriate insulating material will certainly aid in efficient heat transfer. As long as the light irradiation continues, solar absorbers will always be in a higher thermal state than their surroundings. In order to improve the heating efficiencies, various measures can be taken to limit the heat loss. As a basic rule of thumb, a system will have to direct the produced heat to the target with minimal heat resistance, while minimizing the heat transfer to all other areas. Often systems designed for water evaporation and steam generation employ a wicking layer or channels to allow capillary action to draw water instantaneously to the target (surface of the solar absorber).^{19,21,24–31} Since most of these applications take place near room temperature, it would be prudent to design the system such that the solar absorber will transmit the heat

immediately to its specific target, to optimally reduce its operating temperature to directly enhance the photothermal conversion efficiency.

3. Rational material design for enhanced photothermal conversion

Intensive efforts have been devoted to develop high-performance light-to-heat converting materials through a delicate tuning of the composition of solar absorbing materials, design of a multilayer coating and engineering the surface structure in such a way to achieve high solar absorbance with low thermal emittance and optical reflectance.^{10,32,33} Essentially, the key specification that determine the light-to-heat conversion efficiency is the light harvesting ability of the material. Unlike photothermal therapy, which requires a high single wavelength absorption in the near infrared (NIR) range, non-biomedical photothermal solar conversion requires the absorber materials to exhibit highly efficient and broadband optical absorption across the full solar spectrum range (from 250 to 2500 nm) with minimal transmittance and reflectance. Therefore, an overall high solar absorbance of a material is essential to achieve a high solar-to-thermal conversion. In this section, we discuss different strategies to enhance the solar absorption for efficient photothermal conversion using various solar absorber materials, *i.e.* metals, semiconductors and hybrid combinations of organic–inorganic and metallic semiconductors.

3.1. Plasmonic metals

The localized surface plasmon resonance (LSPR) effect of a metallic nanostructure is strongly correlated to the shape, size, dielectric coatings or medium around a particle and the assembly state^{34,35} (Fig. 3). Due to the nature of LSPR, metal-based plasmonic nanoparticles normally absorb at one or a few particular wavelengths, thus limiting their efficiency for solar-thermal conversion.³⁶ To date, gold and silver are the most common materials adopted for photothermal applications. Silver is known for its high plasmonic resonance response and low plasmonic losses in the optical frequency regime, while gold is known for its visible-NIR plasmonic resonances and chemical stability.^{36,37} Recently, other metals have been developed for a wide range of plasmonic applications, such as aluminium, copper, molybdenum, cobalt, nickel and platinum.^{8,36,38,39} In general, reducing the shape symmetry or creating hollow structures could broaden the spectral LSPR band, while changing the size of the particles or the surrounding dielectric will cause mainly a shift in the LSPR band and auxiliary broadening of the absorption to a certain extent.^{9,40}

3.1.1. Influence of the geometry. A pronounced increase in the photothermal heating efficiency due to geometry modifications, such as designing asymmetric structures and elongated, sharp and tapered structures, has not only been predicted by numerical simulations,^{41–44} but also experimentally demonstrated by many research groups.^{20,44} Baffou *et al.* used Green's dyadic function to numerically investigate the mechanisms of the heating effects and

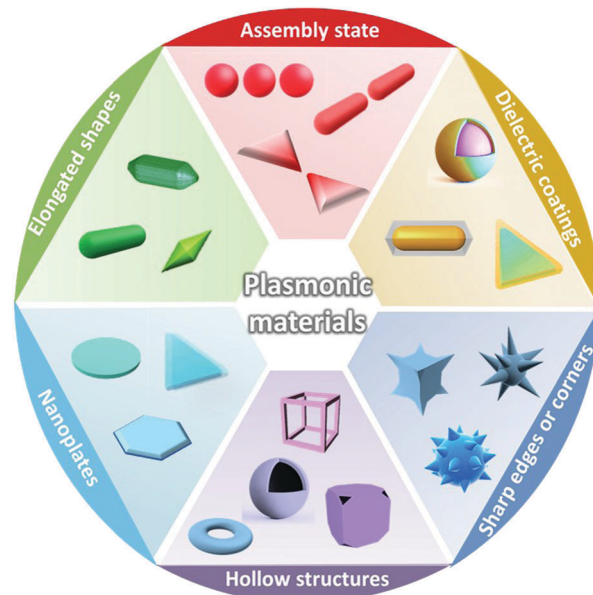


Fig. 3 Various strategies to enhance the optical absorption of plasmonic metals.

efficiencies of homogeneous gold spherical nanoparticles and elongated nanorods colloidal with a constant metal volume in water. This method quantifies the heat generation in gold nanostructures and provides their internal distribution of the heating power density. It was observed that a red-shift occurs when the shape changes from spherical particles to nanorods, together with a substantial increase ($\sim 60\%$) in heating efficiency.⁴⁵ Moreover, plasmonic localized heat generation is often concentrated at sharp tips, corners and edge geometries, where the electromagnetic fields are significantly enhanced.^{9,20} This geometry effect has been validated by many researchers through simulations in assorted sharp geometry structures, including concave structures,^{46,47} nanocubes,^{47–49} octahedrons,⁵⁰ nanoflowers/nanostars,^{32,51,52} nanocages,⁵³ nanoprisms/nanoplates^{54,55} and octopods.⁵⁶ Typically, thin and sharp structures facilitate the incoming electric field to intensely penetrate the constricted geometry, enabling the whole plasmonic matter to be actively involved in heating, while non-symmetrical and hollow structures contribute to broadband absorption. All these geometrical structure factors have explicitly been shown to enhance the overall solar-to-thermal conversion efficiency.

3.1.2. Influence of the dielectric environment. Other than the geometry aspect of the plasmonic material, the optical properties of the plasmonic materials in regard to the absorption resonance wavelength, intensity and range are also sensitive to the surrounding medium. Changing the dielectric environment around the plasmonic particles will result in a shift in the plasmon resonance wavelength.⁵⁷ The surrounding dielectric layer can be made of semiconductors with appropriate bandgaps to enhance the overall light absorption. Chen *et al.* examined the photothermal conversion efficiency of Au nanoparticles with two different coatings, namely Au@ZnS and Au@Ag₂S core-shell structures.¹⁰ Due to the change in the dielectric constant around

the Au nanoparticles, a red-shift in absorption was observed for both the core-shell structures. Even though the absorption intensities of both core-shell structures were the same, the conversion efficiencies of Au@Ag₂S was 1.4 times higher than that of Au@ZnS. This was attributed to the contrasting bandgaps of the Ag₂S (0.9–1.1 eV) and ZnS (3.4–3.8 eV) shells, which resulted in a higher light absorption for Ag₂S at the tested irradiation wavelength (809 nm), while ZnS was transparent to the specific light source. Moreover, it was shown that modification of the dielectric with increasing shell thickness resulted in a large red-shift and a significant enhancement in the plasmon resonance with the Au@Cu₂O core-shell structure.⁵⁸ Furthermore, they also found that the overall extinction spectral was not only determined by the effective shell thickness but also by the porosity of the shell. The increase in shell porosity led to a decrease in the permittivity of the shell, and as such, the less densely packed shell exhibited a narrower absorption spectral range. Hence, it was shown that it was possible to tune the absorption range, intensity and resonance wavelength *via* a dielectric coating of plasmonic materials. Nonetheless, the plasmonic properties may also be adversely attenuated by an inappropriate dielectric coating; hence the constituent, thickness and porosity of the dielectric medium have to be well-controlled for obtaining the desired photothermal effect.

3.1.3. Influence of the intra/interparticle coupling effect.

Plasmonic optical coupling can effectively occur either on an isolated plasmonic nanoparticle (intraparticle) or between two separate neighbouring nanoparticles (interparticle). In order to modify the plasmonic characteristics of an isolated standalone nanostructure, the intraparticle plasmonic coupling scheme is often adopted. Generally, multiple plasmon resonance modes can be accomplished within an isolated nanoparticle, and a new plasmon mode can also be realized. The nanostructure can be configured to host multiple, localized surface plasmon-polariton resonances (SPPR) modes within a compatible interaction range, as such a coupling of individual components of the nanostructure make up the overall optical response. For instance, nanoflowers or nanostars with numerous tips can trigger intraparticle optical coupling between the tip modes and the core modes, which results in a red-shift and augmented dipole moment of the tip plasmons.⁵¹ Moreover, it has been reported that the coupling within the plasmon modes of the closely assembled tips engenders a strong intertip local field enhancement.⁵⁹

It was also demonstrated that the plasmonic coupling effect is significant between adjacent interparticles.⁵¹ When two nanoparticles (*i.e.* a dimer) are placed close to each other, the electric field is concentrated in their gap due to plasmonic coupling. This highly localized field enhancement results in pronounced photothermal effects in the dimer gap. Notably, the overall photothermal performance establishes on both a collective heating ability and a coupling effect between the particles (Fig. 4a). Furthermore, the plasmonic coupling effect often brings about a change in optical absorption, which is highly dependent on the gap distance, the direction of the incident light and the number of particles involved.^{51,60–62}

Hentschel *et al.* fabricated electron-beam patterned planar plasmonic gold hexamers and heptamers to theoretically and experimentally demonstrate that the interparticle gap and the arrangement of the particles play important roles in the plasmonic heating effect (Fig. 4b).⁶³ The total heat generated by two interacting nanoparticles was numerically calculated to be dependent on the polarization of light. When the electric field of incident light is along the axis of the dimer, the total heat generated increases, while when it is in the perpendicular direction of the dimer, the heat generated is diminished.⁶¹ A similar conclusion was drawn by Abajo's group through studying the temperature distribution of dimer particles (Fig. 4c).⁶⁰ The near-field and the temperature distribution corresponding to the dimer configuration are shown in Fig. 4d. It can be observed that the field intensity of the two particles can be tuned by the direction and the wavelength of the incident light, thus they feature different temperature profiles. In addition, the temperature of each sphere as a function of wavelength could be tuned by the gap distance between the dimers (Fig. 4e). Apart from planar interparticle coupling, collective interparticle coupling effect has experimentally been shown to enhance photothermal conversion. For instance, Huang *et al.* synthesized gold nanovesicles consisting of densely packed gold nanoparticles with a specific interparticle orientation. Such interparticle assembling induced an ultrastrong plasmonic coupling effect between the adjacent gold nanoparticles, whereupon the LSPR peak and absorption range were demonstrated to be tuneable by the number of interparticle junctions (Fig. 4f).⁶⁴

In general, plasmon-enhanced electric fields are highly sensitive to the arrangement of the particles and the incident light polarization. However, it is difficult to predict the effect of the photothermal heat generation, especially when the nanoparticles are randomly dispersed in solution. Therefore, in order to take advantage of the coupling effect through the assembly state of the nanoparticles, the arrangement and the design of the particles need to be carefully devised. The possible assembled configurations include highly organized particles on substrates, a hierarchical assembly of nanoparticles, and well encapsulated nanoparticles within a matrix.

3.2. Semiconducting materials

Bandgap engineering governs the light absorption, which in turn controls the photothermal conversion ability of semiconductor materials. Semiconducting materials are generally opaque in the visible and translucent in the IR. Their absorption in the visible spectral region is based on free carrier generation, which depends on the bandgap of the semiconductor, while some transition metal semiconductors have absorption in the NIR region, which has been attributed to either indirect interband transitions or the LSPR effect.^{65,66} The free carrier concentration in semiconductor material is much lower than in metals; however, it can be significantly increased *via* doping or by introducing vacancies, which result in either a shift of the band energy and/or the creation of energy states within the bandgap. The increased in free carrier concentration not only enhances the

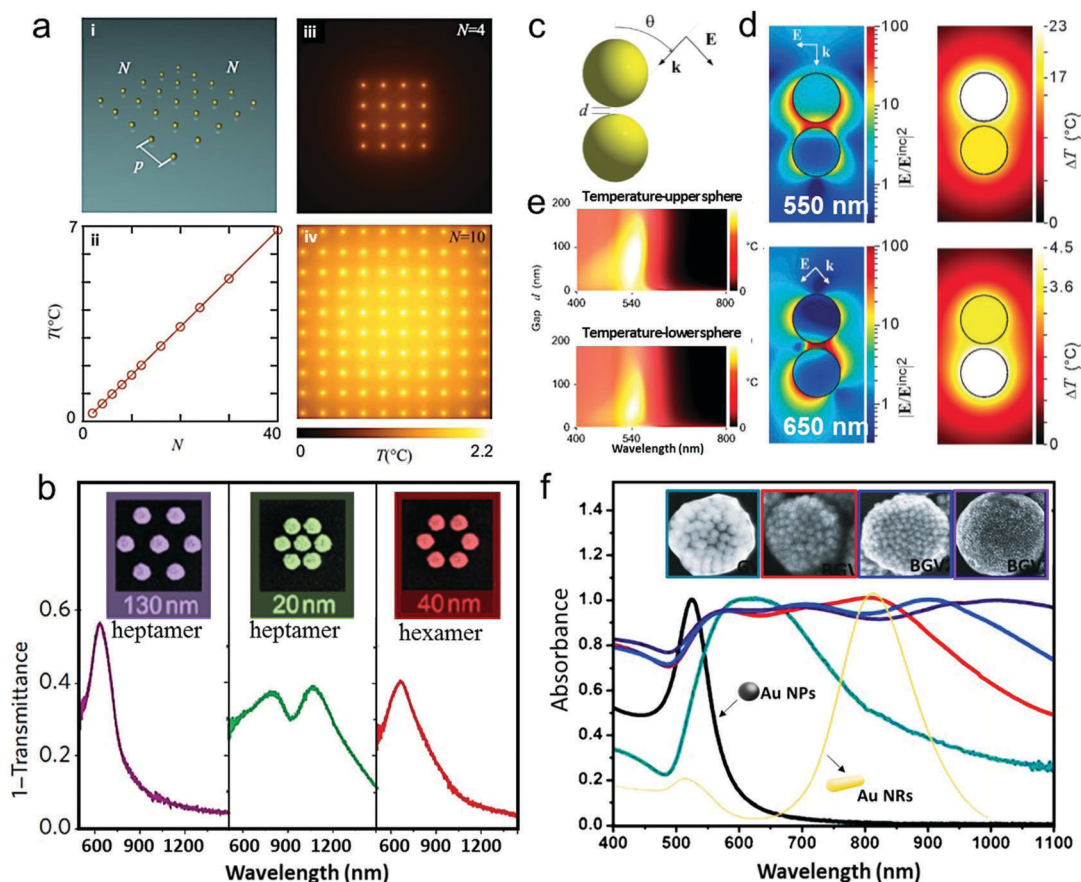


Fig. 4 (a) Numerical simulations of the temperature increase within an $N \times N$ array of gold plasmonic nanoparticles, 40 nm in diameter, $p = 300$ nm in pitch, illuminated from the bottom ($\lambda = 520$ nm, $P = 1$ mW μm^{-2}) showing a 3D model with $N = 5$ (i), the temperature in the centre of the array increases as a function of the size of the array (ii), temperature maps over an array of $N^2 = 4 \times 4$ cells (iii) and $N^2 = 10 \times 10$ cells (iv). Reprinted with permission from ref. 62. (b) Extinction spectra showing the effects of the plasmonic coupling of gold nanodisks with different interparticle gap separations. Reprinted with permission from ref. 63. (c) Sketch of a water-immersed gold dimer with two 100 nm spheres, illuminated by a 1 mW μm^{-2} p-polarized light plane-wave under an angle of incidence $\theta = 45^\circ$ with respect to the axis or rotational symmetry. (d) Distribution of the electric near field (left) and temperature increase (right) under the illumination conditions of $\theta = 0$ and $\lambda = 500$ nm (top) and $\theta = 45^\circ$ and $\lambda = 650$ nm (bottom). (e) Temperature increase in the upper and lower spheres with different gap distances and wavelengths of the incident light. Reprinted with permission from ref. 60. (f) UV/Vis/NIR spectra and scanning electron microscopy (SEM) images (insets) of gold nanovesicles with a controllable interparticle distance with various concentrations in tetrahydrofuran. Reprinted with permission from ref. 64.

absorption of the semiconductor, but also improves the light-to-heat conversion efficiency due to a high probability of nonradiative recombination.

3.2.1. Bandgap engineering. It has been demonstrated that the optical absorption of a semiconductor can be modified as a result of bandgap engineering through adding impurities, self-doping, or introducing disorder to the surface layers to generate donor or acceptor states in the bandgap and/or reduce the bandgap. The smaller bandgap and/or mid-gap energy states are the dominant centres for optical excitation and relaxation, leading to a broadening of the light absorption range of the semiconductor materials.⁶⁷ Ye *et al.* observed a broad absorption in magnesium-reduced TiO_2 nanoparticles populated with oxygen vacancies.⁶⁸ The resulting dark coloured TiO_x exhibited an extended photoresponse from the UV to the NIR spectral regions. Moreover, it was also found that the optical absorption increased with the magnesium concentration,

which brought about a decrease in the oxygen content in TiO_x . Wang *et al.* presented Ti_2O_3 nanoparticles with excellent absorption capability across the whole solar spectrum, showing an average reflectance of less than 10%, lower than that of graphite (20%).⁶⁹ Both the ultrasmall bandgap and nano-sized features of the Ti_2O_3 nanoparticles contributed to the high solar-to-thermal conversion efficiency of $\sim 92\%$. Ti_2O_3 , with a stoichiometric phase of the lower oxidation state Ti_{3p} , can be viewed as an extreme form of oxygen-deficient TiO_2 . When the concentration of oxygen vacancies reaches saturation beyond which the TiO_2 structure is no longer sustainable, the stoichiometric Ti_{3p} phase is then formed. Consequently, oxygen vacancies become the primary defects present in the partially oxidized TiO_2 .⁷⁰ Unlike TiO_2 , which only absorbs in the UV region due to its large bandgap, the bandgap of the Ti_2O_3 semiconductor has been engineered to be much narrower, to absorb the full solar spectrum (Fig. 5a).⁶⁹

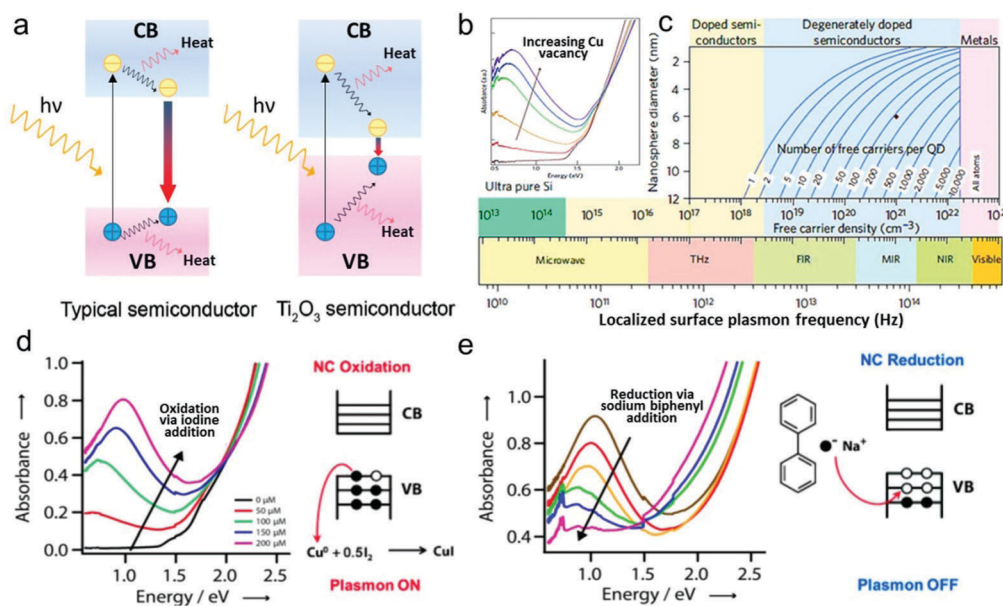


Fig. 5 (a) Illustration of electron–hole generation and relaxation in a typical semiconductor (left) and in narrow bandgap Ti_2O_3 (right). Reprinted with permission from ref. 69. (b) The absorption spectra with increasing Cu vacancy in Cu_2S . (c) LSPR frequency dependence on the free carrier density and doping constraints. Reprinted with permission from ref. 71. (d) The presence of oxidants (iodine) causes the formation of copper vacancies, resulting in holes in the valence band of Cu_2S and an enhanced LSPR. (e) The presence of electron donors (sodium biphenyl) causes the filling of the copper vacancies in Cu_2S , resulting in quenching of the LSPR. Reprinted with permission from ref. 73.

3.2.2. Free-carrier-induced LSPR. Moreover, creating intrinsic defects or introducing extrinsic impurities in transition metal semiconductors has been reported to generate a strong LSPR in the valence band mainly at the NIR region. The absorption characteristics of the semiconductor are similar to those of metals.^{13,66} Hessel *et al.* prepared a Cu_{2-x}Se semiconductor (bandgap 2.1 to 2.3 eV) with a relatively high hole-carrier concentration, which exhibited strong free carrier absorption (540 to 590 nm) due to direct interband transitions. Furthermore, the Cu deficiencies in the nonstoichiometric Cu_{2-x}Se led to a high densities of holes and exhibited an NIR absorbance peak due to surface plasmon resonance.⁶⁶ Similar phenomena have been observed in other semiconductors, such as Cu_{2-x}S nanocrystals,¹⁸ Cu_{2-x}S quantum dots,⁷¹ $\text{W}_{18}\text{O}_{49}$ nanowires¹⁷ and WO_{3-x} nanocrystals.⁷²

Unlike the LSPR of metal, which cannot be dynamically controlled once the nanostructure parameters (size, shape or material) are fixed, the LSPR response of semiconductors can be tuned by the free carrier concentration without modifying the structure. Luther *et al.* demonstrated tunability in the optical properties; as such, NIR absorbance was absent when the copper vacancies in Cu_2S were minimal and conversely, the intensity and blue-shifts in the LSPR band progressively increased with the vacancies (Fig. 5b).⁷¹ Moreover, they showed that the LSPR frequency of the semiconductor nanocrystal could be modulated by the free carrier density and doping constraints (Fig. 5c). By tuning both the size and the doping concentration of semiconductor nanocrystals, the LSPRs could be tuned from the terahertz to NIR frequencies. In order for the LSPR response of a semiconductor to fall in the range of the solar spectra, the required doping concentration must be at

least $\sim 10^{21} \text{ cm}^{-3}$. Alivisatos's group verified this effect through exposing Cu_2S nanorods to both oxidants (iodine) and reducing agents (sodium biphenyl) to facilitate the formation and filling of Cu vacancies, respectively.⁷³ They found that as the amount of oxidant increased, more vacancies were created, which resulted in a gradual blue-shift and an increase in intensity of the LSPR band (Fig. 5d), while as the reducing agent was added, the progressive filling of vacancies resulted in the opposite effect, with a red-shift and reduction in intensity of the LSPR band (Fig. 5e).

3.2.3. Structural design. The other approach to improve the photothermal properties of semiconductor materials is through constructing a hierarchical/superstructure that favours light scattering and that thus enhances photon absorption. For example, Hua *et al.* fabricated hierarchical microstructured HCuPO sheets, which exhibited broad absorption in the vis-NIR band from 460 to 1200 nm together with a high photothermal conversion efficiency of 41.8% at 808 nm.⁷⁴ This was attributed to the d–d transition of Cu^{2+} coupled with a high probability of nonradiative recombination. Similarly, Tian *et al.* synthesized CuS superstructures with faceted end planes of well-defined crystals, which exhibited an NIR photothermal conversion efficiency improvement of 50% as compared to its elemental building blocks (hexagonal nanoplates).⁶⁵ This enhancement was credited to the structural design, whereby the CuS superstructure served as excellent laser-cavity mirrors at 980 nm, which led to superior light absorption ability, thus resulting in an efficient photothermal conversion.

3.3. Hybrid materials design

Another approach to achieve enhanced absorption is through the formation of hybrid materials with complementary and/or

synergetic optical properties due to combinatorial optical responses/characteristics of the composites or favourable structural designs for light scattering. Hybrid materials of various combinations, including bimetallic, metal-organic materials and metal-semiconductor composites, have been developed for enhanced photothermal conversion. Lin *et al.* observed a change in the optical properties of Ag nanocubes with the progressive surface covering of Pt atoms to form Ag-Pt bimetallic nanocages.³³ During the process, the sharp LSPR peak of Ag nanocubes disappeared and was gradually broadened into a wide absorption band, accompanied by the LSPR peak being red-shifted from 510 to 645 nm. This observation was consistent with other metal (Au and Pd) replacements to form bimetallic Ag-Au and Ag-Pd. Besides, hybrid solar absorbers featuring semiconductor nanoparticle@metal-organic frameworks (MOFs) core-shell structures were designed to function as nanoheaters for enhanced photothermal catalysis.⁷⁵ The as-prepared Cu₇S₄@zeolitic imidazolate framework-8 (ZIF-8) hollow core-shell nanoparticles showed a blue-shift of the plasmonic band due to a decrease in the dielectric constant compared with the Cu₇S₄ hollow microsphere core. The cooperative effects between ZIF-8 and Cu₇S₄ components was beneficial for light harvesting. The photothermal transduction efficiency of Cu₇S₄@ZIF-8 was calculated to be 31.1% under the irradiation of a 1450 nm laser.

Optical enhancement can also be achieved through structural design of hybrid materials. In this context, plasmonic particles and 1D columnar scattering media have been designed and demonstrated exceptional light absorption up to 99% for the whole solar spectrum. For instance, Zhu's group combined the plasmonic effect of self-assembled metal

nanoparticles and the wave-guiding effect of aluminium oxide nanoporous nanochannels (Fig. 6a-d).⁷⁶ The nanoporous anodic aluminium oxide (AAO) template functioned as an efficient light trapping medium, which strongly scattered the light through the internal reflections, thus extended the path length of incident light within the nanoporous structure. In addition, the closely packed Au nanoparticles with random sizes and distributions in the aligned pores enabled a high density of hybridized LSPR, which gave rise to broadband absorption. The hybrid material had an impressive average measured light absorbance of 99% throughout the visible to mid-infrared regimes (400 to 10 μm). Similarly, plasmonic metal nanoparticles and vertically aligned microchannels of a wood matrix were exploited to achieve a high solar absorption (99%) across a wide wavelength range of 200 to 2500 nm (Fig. 6e and f).⁷⁷ Natural wood has limited absorption due to the coloured lignin. However, after the incorporation of various plasmonic metals, including Pd, Au and Ag nanoparticles, the colour of the wood turned black and showed enhanced optical absorption in both the range and intensity. Moreover, the absorption was angle-independent owing to the repeated light reflection/scattering and absorption along the unique columnar microstructures of the wood. The efficiency was shown to increase with the light intensity and an efficiency of about 85% under 10 Sun illumination was demonstrated (Fig. 6g). Notably, it is clear that the synergistic effect of combining various solar absorber materials with the preservation of their respective intrinsic optical and/or structural properties can lead to markedly improved optical properties in the absorption range and/or intensity for efficient photothermal conversion.

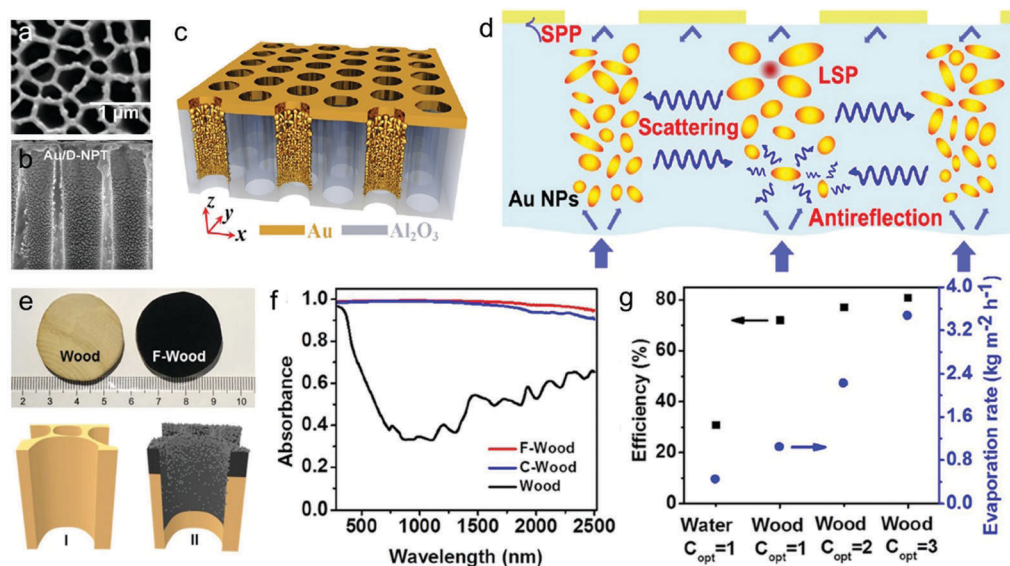


Fig. 6 (a and b) Schematics of plasmonic absorber with Au nanoparticles of random sizes and distributions for the LSPR effect and nanoporous Al₂O₃ templates for effective internal light scattering. (c) 3D schematic of self-assembled Au/AAO plasmonic absorbers. (d) Schematic of an ideal plasmonic absorber. Reprinted with permission from ref. 76. (e) Design of plasmonic wood with both the plasmonic effect of metal nanoparticles and the wave-guiding effect of wood microchannels. (f) Absorption spectra for natural wood and plasmonic wood. (g) Solar-steam-generation efficiency at different light intensities. Reprinted with permission from ref. 77.

4. Photothermal applications

The electromagnetic energy absorbed by solar absorber nano-materials can rapidly be transformed into heat; thus acting as nanosources of heat in the liquid or gas phase, which is a phenomenon that has commonly been manipulated for photothermal vaporization and catalysis purposes. Photothermal vaporization and catalysis can be broadly categorized into volumetric and interfacial systems as well as discrete and collective systems, respectively (Fig. 7). Furthermore, specific types of heat transfers (classical and non-classical) and water contact (direct, indirect, isolation and complete isolation) can be classified according to the implementation of photothermal evaporation (Fig. 7a). Similarly, different types of functions (single and dual heating and catalysis roles) and heat generation (self and assisted heating) can be sorted according to the deployment of photothermal catalysis (Fig. 7b). Thus far, experiments and numerical models have already verified and predicted the thermal distribution of solar absorber surroundings; thereby making it possible to explain the phenomena observed when the heated nanomaterials interact with their immediate surrounding medium/reactant to afford an enhanced evaporation or catalysis performance.

4.1. Photothermal heating vaporization

In nature, solar-driven evaporation is a vital stage in the water cycle, *i.e.* the continuous movement of water between the land,

ocean, rivers, creeks and atmosphere. Only fresh water makes its way to the clouds, leaving behind salt, mineral and metal impurities, constantly refilling the Earth's fresh water supplies, which is essential for life on Earth. However, the natural solar-to-vapour conversion efficiency is limited by optical losses (including poor light absorption and the reflection and transmission of light) and thermal losses (including radiation, convection and conduction losses), which are too low to generate practically useful levels of fresh water. Inspired by this low-cost and clean fresh water generation mode, many efforts have been made to exploit, develop and optimize solar-driven steam generation systems. An efficiency that is often reported is the light-to-steam (or water vapour) generation efficiency. In this case, light (often a solar simulator) is directed at a solar absorber material contained in water. The temperatures at various parts of the systems are also measured to assess the increase in temperature and the uniformity of heating, by using thermocouples or IR cameras. The efficiency η is usually calculated as follows:^{12,78–81}

$$\eta = \frac{m(Lv + Q)}{P_{in}} \quad (4.1)$$

where m is the rate of mass loss of water under illumination, Lv is the latent heat of vaporization of water, normally taken to be 2.26 kJ g^{-1} in the region of interest, Q is the sensible heating of water from an original temperature T_1 to a final temperature T_2 ,

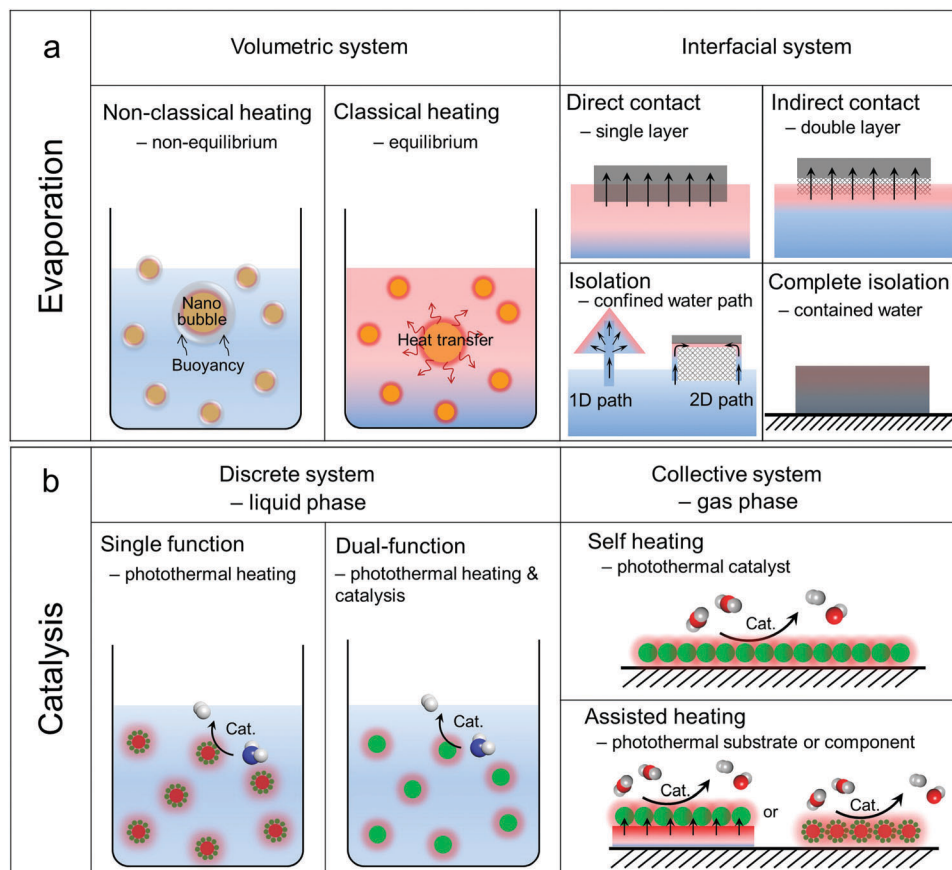


Fig. 7 Schematic diagrams of the different photothermal systems that have been demonstrated for: (a) evaporation and (b) catalysis applications.

given by 2.2, and P_{in} is the incident light power on the solar absorber.

As briefly mentioned, the solar-driven evaporation systems are generally classified into two systems depending on the location of the light-to-heat conversion materials present in the liquid. The first kind involves small particulate absorbers dispersed in a working fluid, which is called a volumetric system, also known as a nanofluid. The second one involves absorbers partially or completely separated from the bulk working fluid, known as an interfacial system (Fig. 8a).

4.1.1. Volumetric system (nanofluid). A volumetric system is a straightforward steam generation route *via* direct nanostructure solar energy absorption. Volumetric absorbers can have surface temperatures lower than the bulk fluid temperature,³⁰ which explicitly reduces significant heat losses to the surrounding media. The steam generation mechanism is still in dispute for volumetric systems. Two plausible mechanisms have been proposed to explain the vapour generation phenomenon.

(a) *Non-classical (non-equilibrium).* The microscopic mechanism of this “nanoparticle-catalyzed” steam formation was proposed by Halas and co-workers.^{12,82} As incident light is irradiated onto Au nanostructure with an overlapping solar spectral (Fig. 8a),

they rapidly heat up and transfer the heat to the immediate surrounding water. The thin shell of water that is in direct contact with the nanoparticles transforms into steam. Since steam is a poor thermal conductor, heat transfer from the heated particles to the water is suppressed. As light continuously irradiates the nanoparticles, the steam shell or “nanobubble” gradually increases until a critical thickness of several hundred nanometres is reached (Fig. 8b). Consequently, the steam/nanoparticle becomes less dense than that of an equivalent volume of water, and so floats to the surface (Fig. 8c). Thereafter, the bubble bursts to release the steam at the surface and the nanoparticle returns to the bulk water. In this way, steam is generated without heating the entire volume of water to the boiling point, thus by-passing energy intensive bulk heating.⁸³ Absorbing nanoparticles suspensions of SiO₂/Au nanoshells and water-soluble carbon black nanoparticles were used by Halas group to demonstrate the particle-nucleated solar steam generation process in an ice bath.¹² Additionally, multiple light scattering between the nanoparticles is also proven to enhance photon absorption, inducing higher temperature. The group reported that 82% of the absorbed sunlight was directly utilized for steam generation, with an overall light-to-steam conversion efficiency of 24%. Moreover, other similar work performed by Halas’s group on

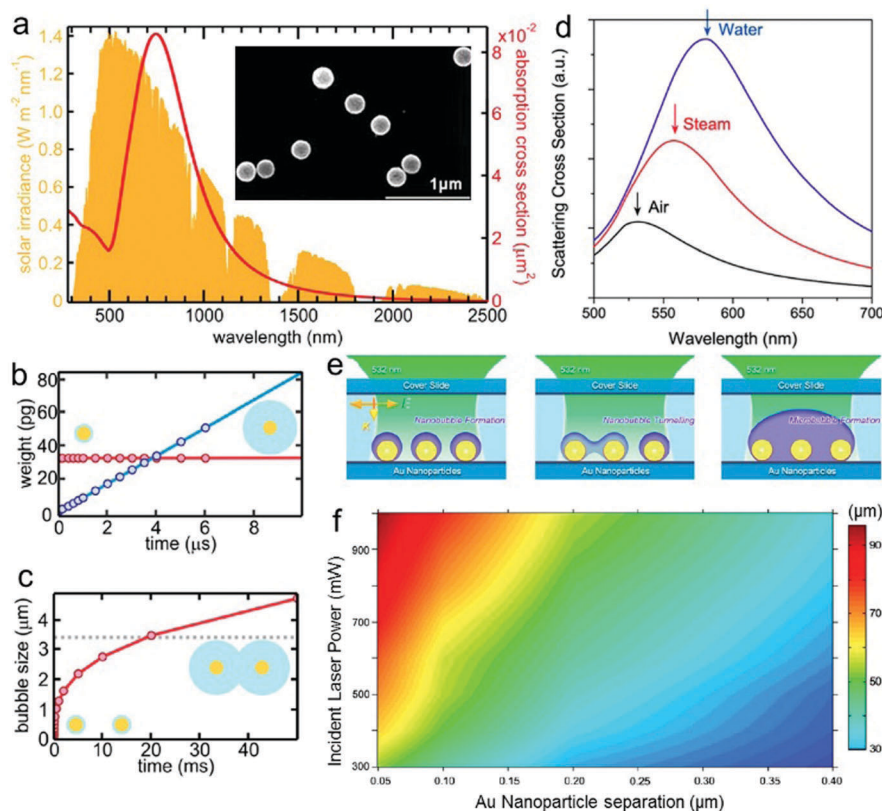


Fig. 8 (a) The absorption cross-section of the Au nanoshells is tuned to overlap the solar spectral irradiance (AM 1.5 G). Inset: SEM image of SiO₂/Au nanoshells. (b) Comparison of the combined weight of a gold nanoshell and its surrounding steam bubble (red curve) with the weight of the displaced water (blue curve) over time (c) size of the steam bubble surrounding a nanoshell over time (red curve). Reprinted with permission from ref. 12. (d) Mie calculation of the scattering cross-sections for a 100 nm diameter Au NP in air (black), surrounded by a 60 nm radius steam bubble (red), and in water (blue). (e) Schematics depicting microbubble formation due to the coalescence of neighbouring nanobubbles of individual Au NPs. (f) Experimentally determined final bubble size and formation time on incident laser power and average nanoparticle separation. Reprinted with permission from ref. 82.

the spectral shift, as calculated from Mie theory, agreed with the experimental scattering spectra using the concentric steam-water shell model (Fig. 8d). Also, the coalescence of nanobubbles into micrometre-sized bubbles was studied. The radii of the nanobubbles were found to increase with the light intensity. When the interparticle spacing was sufficiently small, nanobubbles centred on adjacent nanoparticles touched and coalesced, resulting in the formation of a larger microbubble encompassing several nanoparticles (Fig. 8e). The final size of the microbubble was a function of both the incident power and nanoparticle separation (Fig. 8f). These findings highlight the importance of rationally designing solar-thermal reaction systems, such as the particle concentration, illumination area and reaction container depth, for achieving an efficient solar conversion temperature.⁸⁴

It is important to note that the formation of “nanobubbles” usually requires an intensive laser source ($> 1000 \text{ MW m}^{-2}$) or highly concentrated light ($> 1000 \text{ Sun}$) to overcome the bubble formation intensity threshold.^{12,82,85} Such a high illumination intensity unfavourably requires a high capital investment, including laser generator optical devices and large footprints.^{12,83,84,86,87}

(b) Classical heat transfer (equilibrium). In the classical heat transfer mechanism, well-dispersed solar absorbers absorb the photon energy to produce numerous micro heat sources, which subsequently transfer to the surrounding water and rapidly reach equilibrium with the bulk fluid. In this case, the vapour generation is purely due to the rise in the temperature of the bulk fluid. Chen's group experimentally and theoretically verified the classical global heating process for steam generation under a sunlight concentration of 10 Sun.³⁰ Three water-based nanofluids, namely graphitized carbon black, graphene and carbon black suspensions, were tested as the light receiver. The measured vapour generation efficiencies were approximately the same for all three nanofluids tested ($\sim 69\%$). To elucidate the vapour generation mechanism and validate the experimental results, a consistent set of numerical simulation and analytical modelling runs were performed at the device and nanoparticle scales. The particle model showed that at feasible solar concentrations and illumination intensity, it was highly unlikely to achieve the local temperature gradients that would lead to nanobubble generation around the nanoparticle, as proposed previously. These results clarified that the solar vapour generation of the nanofluids was due to global heating of the bulk fluid, as associated with the classical evaporation phenomena. In addition, other molecular dynamics simulation work also testified that “nanobubbles” under continuous heating conditions cannot be formed even at a high heating power.^{86,88}

For many volumetric solar evaporation systems, steam is generated without heating to the boiling point, indicating light trapping by solar absorber particles that simultaneously absorb and scatter light, resulting in a highly localized heating. Ho's group reported a non-equilibrium steam generation that occurred below the boiling point in the presence of $\text{SiO}_2/\text{Ag@TiO}_2$ core-shell nanoparticles.⁸⁹ The core-shell structures enabled reduced thermal conductivity at the nanoparticle-liquid interface, while functioning as the nucleation sites for

steam generation. The nanocomposites contributed to efficient steam generation through the localized heating effect, as well as by conventional steam generation.

It is worth mentioning that thermodynamic analysis showed that 80% of the absorbed sunlight was converted into water vapour, with inevitably 20% of the absorbed light energy converted into heating the surrounding bulk water.^{12,90} However, these data should be interpreted cautiously, considering the extremely non-uniform temperature distribution within the fluids. Over prolonged light irradiation, the nanoparticle-steam bubble complex progresses and gradually bulk heating arises, which eventually result in simultaneous steam generation below the boiling point and the bulk fluid temperature increasing. Notably, the “nanobubble” hypothesis when combined with the classical heat transfer can provide a more accurate description of the volumetric system after extensive light exposure.^{84,90}

4.1.2. Interfacial system. Volumetric solar-driven steam generation systems, as discussed earlier, typically show a relatively low steam generation efficiency because the bulk working liquid is inevitably heated up too. This makes them susceptible to radiation, conduction and convection heat losses, and also an unnecessary heating of the bulk water that does not participate in steam generation yet consumes considerable energy. To minimize the superfluous energy losses and improve the energy conversion efficiency, an interfacial solar steam generation strategy has recently been explored. In such an approach, both the solar energy harvesting and steam generation are localized at the water-air interface using a non-submersible solar absorber system. In this way, the converted thermal energy is confined at the interface and heats up only the water at the surface. This interfacial confinement reduces heat losses in two ways: first, the temperature of the absorber is lower, thus reducing the radiation and convection heat losses at the absorber surface; second, the temperature of the bulk liquid remains close to ambient temperature, minimizing the heat loss and parasitic use of the converted thermal energy to heat up the bulk water. With this scheme, solar-to-steam conversion efficiencies can reach higher than 90%, as reported in a number of recent studies.^{8,22,26,69,76,91} Thus far, three distinct configurations can be classified in terms of the relative location of the solar absorber to the water surface, as described in the following.

(a) Direct contact configuration. Recently, light absorbing nanoparticles,^{74,81,92,93} films,^{79,80,94-97} monolithic aerogels^{98,99} and foams^{22,100-102} were reported to directly float on water due to their hydrophobic surfaces or lightweight properties, and when illuminated by sunlight, they are capable of efficient steam generation without the necessity of heating the entire fluid volume. Nanoparticles, including plasmonic Au,⁹² carbon particles⁹² and semiconductors,^{74,93} have been investigated. For instance, carbon particles decorated with Fe_3O_4 with a hydrophobic nature and low packing density have shown as high as a 130% increase in saltwater evaporation rate.⁹² These magnetic particles can easily be recycled using a magnet. Similarly, other hydrophobic magnetic nanoparticles, including Fe_3O_4 ,

MnFe_2O_4 , ZnFe_2O_4 and CoFe_2O_4 were self-assembled into floating thin films to demonstrate interfacial solar heating.⁹³ Typically, the water evaporation efficiency was significantly enhanced to ~ 1.7 and 2.2 times higher than that of the nanoparticle suspension and water evaporation itself, respectively.

Besides the freestanding floating nanostructures, self-floating 3D monolithic solar absorber aerogels and foams have been reported.^{22,100} Their inherently porous networks offering water supply passageways, vapour escape channels and heat insulation have seen them widely used in solar steam generation. Hu *et al.* demonstrated free-floating graphene oxide (GO)-based aerogels with tailored properties that displayed highly efficient solar steam generation.⁹⁸ The aerogel consisted of reduced graphene oxide (rGO) sheets, carbon nanotubes (CNT) and sodium alginate (SA), and possessed excellent solar absorption ($\sim 92\%$ across $200\text{--}2500\text{ nm}$), good hydrophilicity and a porous network for heat localization, with 83% efficiency under one Sun illumination (1 kW m^{-2}). Also, a self-floating graphene aerogel synthesized by the photoreduction of GO, exhibited a solar steam generation of ~ 54 and 83% at 1 and 10 kW m^{-2} light intensities, respectively. Other work on rGO synthesized on polyurethane (PU) foam was also demonstrated for an efficient solar steam generation of $\sim 81\%$ at a light density of 10 kW m^{-2} .¹⁰³ Also, hierarchical graphene (h-G) foam consisting of a vertical graphene nanoplates array on a 3D foam skeleton was utilized as a solar absorber (Fig. 9a and b).²² The hierarchical graphene structure achieved high solar spectrum absorption at an arbitrary incident angle with decreased reflection. The external solar-thermal conversion efficiency of the graphene foam reached $\sim 93.4\%$.

Apart from self-floating means, films consisting of an absorber materials' coating on lightweight and hydrophobic scaffolds (stainless steel (SS) mesh, polypropylene (PP) mesh, gauze) have also been reported.^{80,96,97} Wang's research group reported interlaced SS mesh coated with polypyrrole (PPy), which was uniquely showcased as a photothermal membrane with a hydrophobic self-healing capability for interfacial solar heating (Fig. 9c).⁹⁴ PPy was chosen as the photothermal polymer owing to its broad spectrum absorption and its self-healing/restoring hydrophobicity property. The water evaporation rate of a PPy coating thickness of $2.6\text{ }\mu\text{m}$ was shown to be $0.92\text{ kg m}^{-2}\text{ h}^{-1}$, equivalent to a conversion efficiency of $\sim 58\%$. Likewise, Jiang's group showed that a carbon black coating on a superhydrophobic gauze film was able to float on the surface of water and facilitated interfacial heating.⁸⁰ Control experiments were conducted under different light intensities and proved that the floating black gauze could achieve an evaporation rate $2\text{--}3$ times higher than that from natural evaporation. Other interlaced PP and SS meshes coated with black TiO_x nanoparticles were also reported to be superior light absorbers for interfacial water evaporation.^{68,95}

(b) *Indirect contact configuration.* Generally for the indirect contact configuration, bilayered structures are designed such that the top layer is a solar absorber material that absorbs the solar flux without contacting the bulk water, while the bottom layer limits conduction of the generated heat to the underlying water body. Additionally, the bottom layer offers a mechanically stable support for the solar absorber to prevent material dispersion into the water body. Airlaid paper,^{104–106}

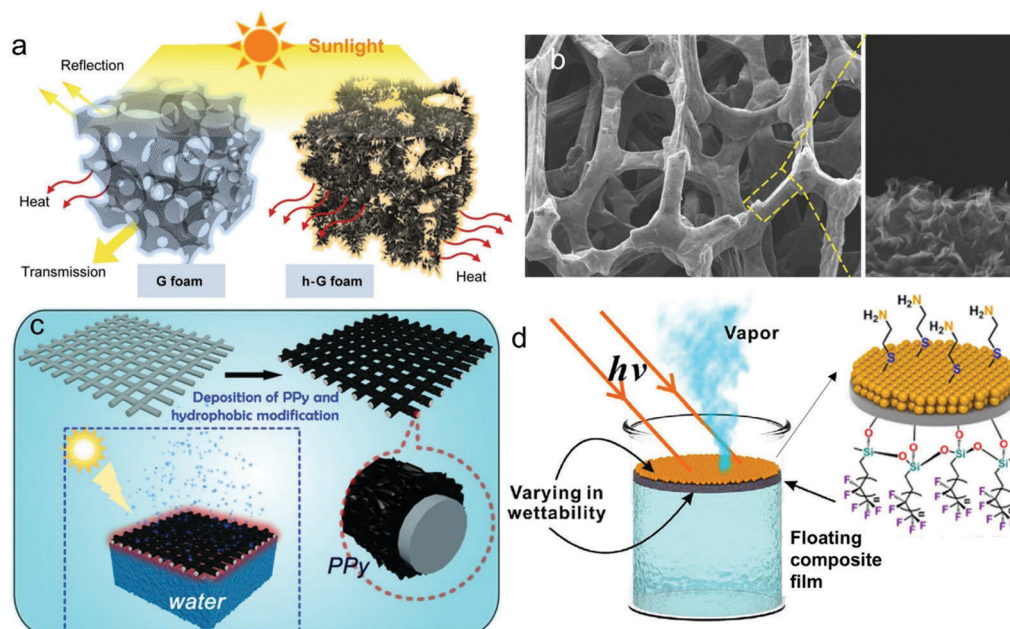


Fig. 9 (a) Schematic of solar-thermal conversion difference between ordinary graphene foam and h-G foam. (b) SEM image of the h-G foam in low and high magnifications. Reprinted with permission from ref. 22. (c) Schematic illustration of the preparation of the self-healing hydrophobicity photothermal membrane (PPy-coated SS mesh) for interfacial heating. Reprinted with permission from ref. 94. (d) Schematics depicting floating double-layered film (top: light-to-heat plasmonic conversion layer; bottom: supporting AAO membrane) with varying wettability. Reprinted with permission from ref. 108.

AAO membrane,^{8,76,107,108} carbon foam,⁷⁸ biofoam,²³ wood,^{27,109,110} silica,¹¹¹ cellulose^{69,91,112,113} and polyvinylidene fluoride (PVDF)¹¹⁴ membrane have all been employed as the backing or lining layer due to their porous structure, low density and heat insulating properties. Airlaid paper, a textile-like porous hydrophilic material with low thermal conductivity ($0.03\text{--}0.05\text{ W m}^{-1}\text{ K}^{-1}$), is a commonly used thermal insulator support.¹⁰⁶ It possesses a high surface roughness which can produce multi-scattering of incident solar light, yielding enhanced photon absorption. The innumerable microscale pores of the paper can be leveraged for capillary water flow to the hot zone, leading to the rapid replenishment of water. Furthermore, its low thermal conductivity acts as a good thermal insulator to prevent heat transfer from the solar absorber film to the bulk water. Such an increased light absorption and reduced heat loss enable airlaid-paper-based gold nanoparticle (AuNP) film to exhibit a much higher evaporation efficiency (77.8%) than a freestanding AuNP film (47.8%). Based on the same function as the airlaid paper, cellulose and PVDF membranes were also used as lining platforms for coating solar absorber materials ranging from rGO,^{105,112} zeolite,¹¹³ Ti_2O_3 nanoparticles,⁶⁹ plasmonic Au nanoparticles¹¹⁵ to Ti_3C_2 MXene,¹¹⁴ all of which demonstrated enhanced interfacial solar evaporation. Among these, the highest efficiency could reach up to 90% under 3 Sun.⁹¹

The foam-like structure of an outstanding thermal insulation has been judiciously used as a thermal insulator support. Chen's group developed a double-layer structure consisting of a carbon foam support and an exfoliated graphite layer. The bottom carbon foam was thermally insulating with a smaller pore size for liquid supply, while the top exfoliated graphite layer had a larger pore size for the vapour to escape. Solar-thermal efficiencies up to 85% at 10 kW m^{-2} and 64% at 1 kW m^{-2} were reported.⁷⁸ Other bilayered foams in the form of a hybrid biofoam of rGO and bacterial nanocellulose also exhibited excellent solar steam generation of 83% at 10 kW m^{-2} .²³

An AAO nanoporous support was also exploited in a bilayered solar evaporation scheme. The AAO nanoporous template presented a reduced refractive index and a small impedance mismatch that enabled an extremely low reflectance and efficient light coupling within the columnar pores, therefore significantly enhancing the light absorption. The AAO bilayer schemes, e.g. AuNP/AAO and Al NP/AAO, were reported to be able to achieve over 90% efficiency under concentrated Sun.^{8,76} Moreover, Yu *et al.* fabricated a bilayered evaporation structure with the top light-to-heat conversion AuNP film and the bottom supporting AAO membrane (Fig. 9d).¹⁰⁸ The results demonstrated that the evaporation rate could be tuned by chemically modifying the AAO-based AuNP film with different wettability, meaning it could behave like a "water gate" that controls the rate of water vapour flow.

A silicone-based porous material with low thermal conductivity and poor light harvesting was also deployed as a good thermal insulator layer. A bilayered structure, consisting of a top self-floating hydrophobic CNT membrane and a bottom hydrophilic macroporous silica support was used for solar evaporation.¹¹¹ The silica served triple functions: (a) as a water

pump that continuously wicks water through its interconnected macro-sized channels to the top heating layer *via* capillary action, (b) as a heat transfer barrier that reduces the generated heat from being transferred to the bulk water, and (c) as a mechanical backing that supports the top thin CNT photo-thermal layer and at the same time realizes a seamless connection between two distinctive layers. A water evaporation rate as high as $1.32\text{ kg m}^{-2}\text{ h}^{-1}$ with a solar-thermal conversion efficiency up to 82% was achieved.

Apart from man-made structures, natural wood with vertically aligned microchannels serves to deliver water, ions, and other nutrients during growth and can be exploited as a solar absorber and support for water evaporation. A bilayered structure composed of radially cut wood and a top coating of GO showed a solar-thermal efficiency of $\sim 83\%$ under a power density of 12 kW m^{-2} .¹⁰⁹ Owing to its abundance, biocompatibility, hydrophilicity, low thermal conductivity and high optical absorption, wood with a natural vessel structure makes an excellent supporting layer candidate for solar steam generation. Other works on wood used for interfacial solar evaporation have also been reported.^{27,110}

(c) *Isolation configuration.* Despite tremendous progress in interfacial water evaporation in the past few years, heat loss between the hot interfacial medium and bulk water is inevitable due to the large contact area. To further minimize heat conduction losses, a variety of structures that separate the solar absorber from water bodies have been designed and fabricated to improve the energy conversion efficiency.²¹ Confined 1D or 2D paths are constructed to link up the solar absorber and water body for efficient water supply. It is worth mentioning that nature already provides an elegant solution for efficient evaporation. In the transpiration processes of plants, water is pumped from the roots and transported from the bottom to the top through xylem vessels and lumina (channels) inside the trunks and branches, before finally being released to the atmosphere through the leaves to achieve an efficient water supply and evaporation. In addition, the natural 3D structures of plants maximize light absorption from a wide range of incident angles. Inspired by plants, a 3D artificial transpiration device, composed of a 3D hollow cone absorber connected with a 1D water path has been reported.³¹ With all three components of heat loss and the angular dependence of light absorption minimized, the transpiration device displayed a solar steam efficiency over 85% under one Sun. Furthermore, a similar conceptual idea of a confined 1D water pathway connected to a solar absorber was demonstrated using a living organism *fungi mushroom*. It was validated as an efficient solar steam-generation medium. The umbrella-shaped, porous structure, and fibrous stipe of the mushrooms not only provided efficient light absorption, water supply, and vapour escape, but also simultaneously suppressed the three components of heat losses (Fig. 10a and b).²¹ The carbonized mushrooms achieved $\sim 62\%$ and $\sim 78\%$ conversion efficiencies under 1 Sun illumination, respectively (Fig. 10c and d).

Besides the umbrella-shaped structure, a flat film equipped with a wicking fabric and insulating foam was also reported.

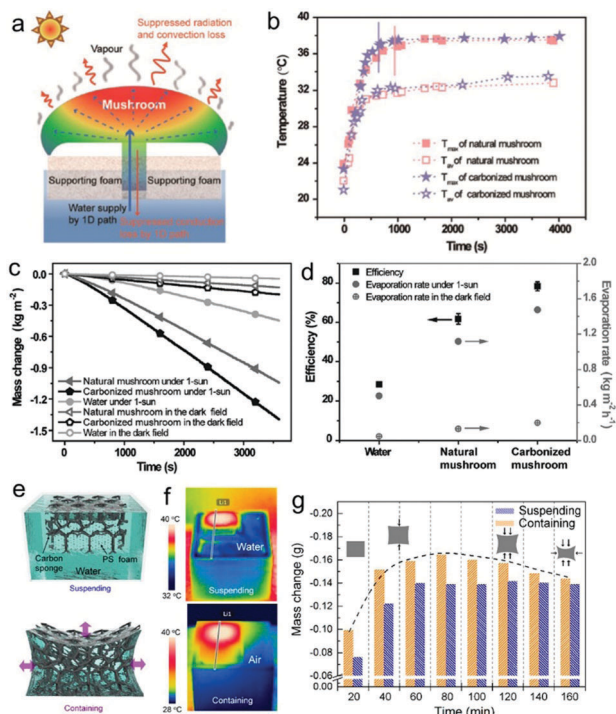


Fig. 10 (a) Schematic illustration of the heat behaviour in a mushroom-based structure. (b) Maximum and average temperatures of surfaces for natural and carbonized mushrooms as a function of time. The blue and pink lines show the turning points to the quasi-steady state of natural and carbonized mushrooms. (c) Mass change of water over time in the dark field and under 1 Sun illumination. (d) Solar steam efficiency (black, left-hand side axis) and evaporation rate in the dark field and under 1 Sun illumination (grey, right-hand side axis) with the natural and carbonized mushrooms. Reprinted with permission from ref. 21. (e) Schematic diagrams of CS in suspending and containing configurations. (f) The corresponding IR images of the suspending and containing configurations after 30 min light irradiation. (g) The evaporation mass loss comparison of CS in contact and completely isolated from bulk water. The insets show the self-contained CS volume change with evaporation time. Reprinted with permission from ref. 116.

Chen's group demonstrated a floating solar receiver capable of generating 100 °C steam under ambient air conditions without optical concentration.²⁵ A variety of low-cost and commercially available materials, such as cermet-coated copper substrate, polystyrene foam and bubble wrap, were used to construct the solar receiver system. By varying the thermal concentration, the receiver could generate saturated steam at 100 °C, or low-temperature vapour at high efficiencies (64%). A similar architecture using a long-range vertically aligned graphene sheets membrane was reported as an efficient solar-thermal converter along with a polystyrene foam heat insulator and glass fibre water wicker.²⁶ Water evaporation rates of 1.62 and 6.25 kg m⁻² h⁻¹ under 1 and 4 Sun, equivalent to outstanding efficiencies of 86.5% and 94.2%, were obtained, respectively.

A 2D water path, *e.g.* hydrophilic porous paper, could also realize efficient water supply in isolated evaporation configuration. For instance, Zhu reported a polystyrene foam (thermal insulator) that was wrapped with porous cellulose (for a 2D water path) with GO film (solar absorber) on the

top surface.¹⁹ By doing so, the isolated solar desalination system displayed efficient 2D cellulose water supply and suppressed the thermal losses, which enabled an efficient (80% under 1 Sun illumination) and effective (four orders salinity decrement) solar desalination. Another similar work was reported on carbon black coated paper affixed with expanded polystyrene foam with the edges of the paper in contact with bulk water.²⁸ Due to the thermal isolation between the surface liquid and the bulk water, a thermal efficiency of ~88% was obtained under 1 Sun, corresponding to an evaporation rate of 1.28 kg m⁻² h⁻¹.

Ho's group reported a novel elastic cellular sponge that self-confined and stored water and as such it did not need to stay in water.¹¹⁶ The ultralightweight nitrogen-enriched carbon sponge, a 3D elastic cellular solid, could soak up water and perform efficient *in situ* solar evaporation. The carbon sponge possessed a favourable inbuilt structural hierarchy with mesoporous fibres that were seamlessly interconnected to form elastic macroporous open cells. The sponge capillary action was exploited to wick and confine water so as to completely isolate it from the bulk water to perform evaporation remotely with a high efficiency of 90%, which was higher than any common suspending configuration (Fig. 10e–g).

4.2. Photothermal catalysis

Photocatalysis has been widely investigated for the direct conversion of solar energy to chemical energy and for sustainable environmental remediation in the fields of water photolysis,^{14,117,118} CO₂ conversion,^{119,120} organic photosynthesis^{121,122} and pollutant degradation.^{123–126} The definitive advantages of photothermal catalysis over traditional thermal- or photocatalysis are: (1) a more efficient utilization of the full solar spectrum; (2) no input of thermal energy for the catalysis processes; (3) the temperature on the catalyst surface rises instantaneously and the heat is confined on the surface of the catalyst (localized heating effect); (4) in some cases, photothermal catalysis can effectively refrain the deactivation of the catalyst and increase the selectivity for the desired products.¹⁵ Hence, such an analogous photothermal-assisted catalysis is a highly feasible and versatile strategy, without having to impose high requirements on the surface properties of the catalyst.^{16,75} So far, two photothermal heating configurations of discrete and collective systems have been widely used in liquid- and gas-phase catalytic reactions, respectively (Fig. 7b).

4.2.1. Discrete system. Generally, many catalytic reactions are carried out in a liquid medium, such as hydrogen production,^{127–129} hydrogenation,^{50,130,131} aromatic alcohols oxidation¹³² and 4-nitrophenol reduction.^{133–138} Photothermal catalysts in these cases are often dispersed uniformly in solutions, which are similar to the solar absorber in a volumetric system for photothermal evaporation. Photothermal catalysts absorb the photon energy to serve as nano heat sources at the catalyst/solution interface, which show a higher surface temperature, leading to a higher catalytic activity. Subsequently, the generated heat transfers to the surrounding solution, thus also resulting in a substantial global temperature rise.

(a) *Single function.* Similar to the indirect heating configuration, a combination of two different materials with respective catalytic and photothermal properties can be integrated into a hierarchical structure to yield a more efficient catalytic performance in liquid-phase reactions. In such a hierarchical structure, the heat generation source is deliberately placed near to the photocatalytic active sites, in order to efficiently deliver heat for boosting chemical reactions.⁹ Xu's group recently reported accelerated cyclocondensation reactions based on the hierarchical core-shell $\text{Cu}_7\text{S}_4@\text{ZIF-8}$ nanostructure, where the core and shell of the composite had photothermal and catalytic functionalities, respectively (Fig. 11a).⁷⁵ The hierarchical Cu_7S_4 hollow microsphere core acted as a plasmonic nano-heater with a photothermal transduction efficiency of 31.1% and the surface temperature increased up to 94.0 °C under the irradiation of a 1450 nm laser (Fig. 11b). The localized heat directly acted on the surrounding ZIF-8 shells with synergistic acid-base catalytic sites, which led to a 4.5–5.4 fold improvement in catalytic cyclocondensation reactivity compared to that of the process at room temperature (Fig. 11c–e). Another illustration was made on Ni/rGO and Cu/rGO nanocomposites for the reduction of 4-nitrophenol to 4-aminophenol with NaBH_4 .^{137,138} Under NIR irradiation, the reduction rate was significantly enhanced due to the excellent NIR photothermal conversion property of rGO, which induced heating of the local environment near the Ni or Cu nanoparticles and the surrounding reaction medium.

(b) *Dual function.* Instead of a single function that separately delivers heating and catalysis roles, dual-functional photothermal conversion and catalytic reactive nanomaterials have

been reported.^{50,130,132} For instance, noble metals and their alloys are widely used for concurrent surface plasmon-driven photothermal conversion and catalysis reactions. Pd concave nanostructures with a high cross-section and reduced shape symmetry were synthesized to enable broad light spectral harvesting for photothermal enhanced hydrogenation.⁵⁰ The catalytic efficiency under 100 mW cm^{-2} full-spectrum irradiation at room temperature turned out to be comparable to that of thermally (70 °C) driven reactions. Moreover, the Pd concave nanostructures possessed a large quantity of atoms at the corners and edges, where local heat was more efficiently generated to simultaneously provide active sites for the reaction. Furthermore, Pd in MOFs has also shown its dual function in photothermal conversion and catalysis for the hydrogenation of olefins.¹³¹ Remarkably, the catalytic efficiency under 60 mW cm^{-2} full-spectrum or 100 mW cm^{-2} visible light irradiation at room temperature was comparable to that of a process driven by heating at 50 °C. Besides Pd, Au nanoparticles have also displayed plasmonic and catalytic properties to transduce light to heat a surrounding bulk solution to a temperature of 65 °C, resulting in a higher reduction of 4-nitrophenol. Likewise, Au^{133,134,139} and the alloy AuPd^{130,136} have been reported for localized heating and catalytic reactions.

Alternatively, a noble metals combined with other photothermal or catalytic materials have also shown enhanced photothermal catalytic activity. For instance, Pt/PCN-224(Zn), which consists of Pt nanocrystals and the porphyrinic MOF PCN-224(Zn) exhibited excellent catalytic performance at ambient temperature for aromatic alcohols oxidation.¹³² Similarly, Song *et al.* used Pt/TiO₂ to investigate its dual photothermal heating

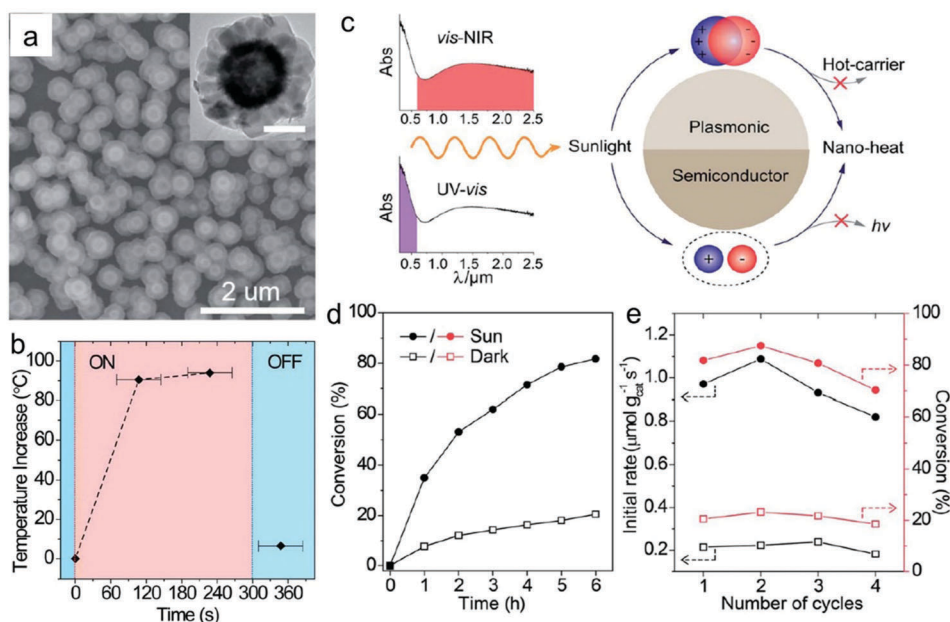


Fig. 11 (a) SEM and transmission electron microscopy (TEM) (inset) images of the core-shell $\text{Cu}_7\text{S}_4@\text{ZIF-8}$ nanoparticles. (b) On/off response of the surface temperature increase of Cu_7S_4 nanoparticles. (c) Schematic depiction of the full-solar-spectrum photothermal effect of plasmonic semiconductor nanomaterials via plasmon- and exciton-based approaches. (d) Catalytic performance of $\text{Cu}_7\text{S}_4@\text{ZIF-8}$ nanoparticles with or without simulated sunlight illumination. (e) Recyclability test (1st hour rate and 6 hour conversion) of the core-shell photothermal catalyst. Reaction conditions: 100 mW cm^{-2} full-spectrum irradiation, 11.2 cm^2 illumination area, 25 mL reactor, room temperature. Reprinted with permission from ref. 75.

and catalytic hydrogen generation by formic acid reforming.^{137,138} At 90 °C, the photothermal reaction rate was shown to be 8.1 and 4.2 times that under photo or thermal conditions alone, indicating photo- and thermal-synergetic coupling effects.

4.2.2. Collective system. Different from the liquid-phase catalysis, the gas-phase catalytic reaction is often established on catalyst loading on a support, *e.g.* glass slide,³ quartz film¹⁴⁰ and quartz glass,¹⁴¹ which is then placed in a gas reactor. Under light irradiation, the heating-up of an ensemble of nanostructured catalysts manifests in a collective heating process, resulting in a dramatic increase in the temperature of the catalyst. Due to the absence of liquid, the temperature of the catalyst can reach tens to hundreds of degrees Celsius. The generated high temperature is adequate for overcoming the energy barrier to initiate the catalytic reaction, while the lower temperature can also sufficiently accelerate the reaction rate. Photothermal catalytic reactions for CO₂ reduction, CO oxidation and volatile organic compounds (VOCs) degradation using such a collective heating configuration have been widely tested and proven.

(a) Self heating. A series of catalysts, including ultrathin Mg–Al layered double hydroxides (LDHs) matrix/Ru nanoparticles (Ru@FL-LDHs),¹⁴² Group VIII metals (Ru, Rh, Ni, Co, Pd, Pt, Ir and Fe),⁵ elemental boron¹⁴⁰ and Fe@C core-shell hybrids,¹⁴³ have achieved efficient photothermal catalysis under collective photothermal heating systems, as reported by Ye's group. The performance of Group VIII metals in photothermal catalytic CO₂ conversion was first investigated.⁵ Supported by the inert Al₂O₃, all of the Group VIII catalysts were shown to be excellent solar absorbers over a very broad spectrum range that covered UV, Vis and IR radiation. As such, they could maintain a high reaction temperature of 300–400 °C under irradiation for efficient photothermal CO₂ conversion. Ru/Al₂O₃ exhibited the highest reaction rates, followed by the Ni/Al₂O₃ and Co/Al₂O₃ and Pd/Al₂O₃ catalysts. The highest CO₂ conversion of 96.26%, selectivity for CH₄ of 99.51% and selectivity for CO of 95.96% were achieved by Rh-, Co- and Fe-supported Al₂O₃, respectively. Moreover, monochromatic visible light filters were used to suppress the thermal effect (*ca.* 21–23 °C), which showed insignificant CH₄ evolution, suggesting a mediation of the photothermal effect rather than the contribution of a photocatalysis reaction. Another efficient photothermal catalyst for CO₂ reduction is elemental boron (Fig. 12a and b).¹⁴⁰ In the presence of only water, the continuous formation of CO and CH₄ was realized with high reaction rates of 1.0 and 2.5 mmol h⁻¹ under full arc irradiation, respectively. Notably, the amorphous boron catalyst was still active under visible light irradiation with a fairly high production of CO at 0.8 mmol h⁻¹ and CH₄ at 1.9 mmol h⁻¹. It was proposed in this process that the boron particles were photothermally heated to a high temperature (378 °C, Fig. 12c), which enabled the localized hydrolysis of boron with water, leading to the *in situ* generation of H₂ and boron oxides. Then, the reduction of CO₂ to CO and CH₄ was photothermo-catalyzed by the illuminated boron particles with the *in situ* formed H₂ as the proton source and electron donor (Fig. 12d–f).

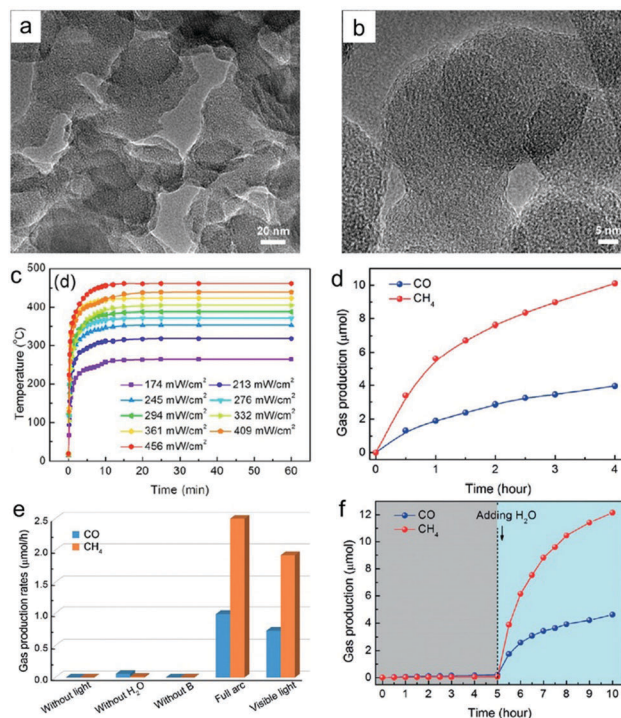


Fig. 12 (a) TEM and (b) high-resolution TEM images of the amorphous boron material. (c) The increase in temperature of amorphous boron samples under light irradiation at different light intensities. (d) Gas evolution (CO and CH₄) over the amorphous boron catalyst under full arc irradiation with a Xe lamp. (e) Gas-evolution rates under different conditions. (f) Time course of CO and CH₄ generation upon irradiation of the amorphous boron catalyst before and after the addition of water to the reaction system. Reprinted with permission from ref. 140.

Furthermore, Li's group reported a solar-light-driven thermocatalysis mechanism for the oxidation reaction of CO and VOCs, *i.e.* benzene, toluene and acetone.^{3,6} When the temperature of the photothermal manganese oxide mineral catalyst increased above the temperature at which the catalytic reaction starts, the thermocatalytic oxidation reaction kicked in. It was experimentally verified that no conventional photocatalysis occurred on the catalysts, while thermocatalytic oxidation following the Mars–van Krevelen mechanism took place instead.^{144,145} Here, basically, organic molecules adsorbed on the surface of the catalyst are oxidized by lattice oxygen, and the resultant oxygen vacancies are subsequently replenished by O₂ gas-phase. With the photothermal activation of the layered manganese oxide mineral catalysts, considerable improvement in the CO degradation was achieved due to significant production of lattice oxygen upon solar light irradiation.

However, in many cases, when the photothermal-induced collective temperature is not high enough to provide adequate energy to overcome the successive barrier, other initiation means, *i.e.* photo or conventional thermal energy, can also enhance the catalysis reaction by augmenting the migration and diffusion of reactants and products. Zhao's group demonstrated a synergistic photo/thermocatalytic effect for CO₂ reduction based on an oxygen-deficient mesoporous WO₃ (m-WO_{3-x}) catalyst.¹⁴⁶

With the photothermal effect, the CH_4 yield showed an additional $\sim 20\%$ improvement. The mechanism was proposed whereby the temperature on the catalyst surface rises quickly, which augmented electron excitation and relaxation, leading to intensified photocatalysis. The photogenerated holes of m-WO_{3-x} are trapped by oxygen ions generated from CO_2 , which is beneficial for the separation of photogenerated carriers and which consequently improves the photocatalytic efficiency. The same mechanism was described in double perovskites $\text{LaSrCoFeO}_{6-\delta}$ (LSCF) and $\text{LaSrCoFeO}_{6-\delta}$ with a 3D ordered macroporous structure (3DOM-LSCF) catalysts for CO_2 reduction.¹⁴⁷

(b) *Assisted heating.* Another route to improve the catalytic performance by photothermal-assisted heating is derived from a coupled solar absorber component or substrate. Conveniently, utilization of the nanostructured substrate with photothermal properties provides a collective heating source for enhanced light absorption and heat conversion, representing an effective strategy for improving the photocatalytic conversion rates of gas-phase reactions. Ozin's group developed a "black silicon" support for photocatalysts dispersion to enhance the hydrogenation of CO_2 reactions. The "black silicon" was composed of vertically aligned Si nanowires (SiNWs) etched into a Si wafer and exhibited a reflectance less than $\sim 3\%$ throughout most of the solar spectrum.^{148,149} Ru^{150} and $\text{In}_2\text{O}_{3-x}(\text{OH})_y$ ¹⁵¹ nanoparticles were then evenly coated on Si NWs to carry out CO_2 reduction. In the case of $\text{In}_2\text{O}_{3-x}(\text{OH})_y/\text{SiNW}$, solar photons with an energy greater than the bandgap of $\text{In}_2\text{O}_{3-x}(\text{OH})_y$ (> 2.9 eV) contribute to the photocatalytic reduction of CO_2 , while sub-bandgap photons (< 2.9 eV) are absorbed or trapped by the SiNW arrays, generating heat to drive the reaction using only light energy without external heating. The $\text{In}_2\text{O}_{3-x}(\text{OH})_y/\text{SiNW}$ photocatalytic CO_2 reduction rates could be significantly increased to ~ 6 -fold over the pristine $\text{In}_2\text{O}_{3-x}(\text{OH})_y$, attributed to the reduced reflection losses and improved light harvesting efficiency. Moreover, another combination of assisted heating involving two functional composite materials with photothermal and catalytic properties has proven to be an effective pathway for introducing the light-induced heating effect into catalysis. Au is not an efficient thermal catalyst for oxidizing ethanol to CO_2 . However, there exists a synergism between Au and UV-illuminated TiO_2 in the photothermal regime, with over 50% and 100% increases in catalytic performance seen in comparison to heating TiO_2 under visible light and UV illumination, respectively.¹⁵² Photothermal catalytic results and detailed probing of the postreaction surface carbon species on Au/TiO_2 indicated that the photo enhancement under UV illumination was due to the congruent roles of the photo and thermal catalyses. The photo enhancement under visible light illumination was ascribed to the plasmonic-mediated electron charge transfer from the Au to the TiO_2 support. Also, the electron charge transfer promoted C–C bond cleavage by the Au deposits.

5. Conclusions and perspectives

In summary, herein we presented a detailed review of diverse solar absorbers as well as covering the materials and system

design strategies that are widely used in photothermal vaporization and catalysis. To date, various class of materials been developed with high solar absorptance for efficient light-to-heat conversion, including noble metals, semiconductors, carbon and a combination of organic–metallic and organic–semiconductors. These solar absorbers have been shown to be effective nanoheaters that can achieve highly localized and selective thermal spots for enhanced vapour/steam generation and catalytic fuel production. By and large, these materials are devised to possess nanoscale dimensions, hybrid material constituents, beneficial geometrical features, a porous framework and a hierarchical structure in view of establishing one or a few virtues of a shorter transfer length, higher surface area, induced plasmonic effect, narrowing of the bandgap and modulation of the carrier density *etc.* to effectively enhance the photothermal performance. Moreover, the coupling of a low thermal conductivity insulator to solar absorber materials helps to prevent excessive heat transfer/loss from the solar absorber film to the bulk fluid or ambient environment, which further augments the photothermal conversion efficiency.

In spite of the recent achievements in high photothermal conversion efficiency, there is still a lack of *in situ* and *operando* (experimental and computational) studies conducted to clearly elucidate the photothermal mechanism of various processes, especially complex photothermo-catalytic transformation. Also, more efforts can be dedicated to improve and/or construct rationale experimental set-ups, *operando* spectroscopy and microscopy to accurately determine and image local temperature, thermal conductivity, photon-to-thermal energy conversion, heat transfer/loss, *etc.* at a macroscopic and nanoscopic level. Moreover, it is important that near-field and far-field techniques that are robust and that do not suffer from either complex calibration procedures or artefacts (sensitive to factors such as thermally induced diffusion or excitation intensity fluctuation) are made accessible for in-depth photothermal characterization and understanding. Finally, further advancements in photothermal vaporization and catalysis capabilities not only demand the synthesis of low-cost, well-controlled nanostructure but also those that are endowed with hybrid photonic, thermionic and electronic functionalities that can be spatially modulated. Moreover, advanced nanomaterials should be judiciously tuned and tailored to bring out the unique electronic and optical properties to display the localized surface plasmonic resonance, quantum confinement, *etc.* phenomena accompanied by a large surface area, high surface energy and ideal geometrical structure, all of which are highly favourable for the photothermal effect. Finally, further advancements in photothermal vaporization and catalysis capabilities not only demand the synthesis of low-cost, well-controlled nanostructure but also those that are endowed with hybrid photonic, thermionic and electronic functionalities that can be spatially modulated. We believe that more highly efficient and sustainable photothermal materials would be developed in the future not only for water evaporation and catalysis applications but also for use in photothermal ablation, therapy, sensor and electricity generation fields.

Conflicts of interest

There are no conflicts to declare.

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